



The Phase Problem: Examples

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Kinematical




- Easy to interpret images
- Diffraction data can be inverted

When is image interpretation easy?




- Frequently
- Obscurations of dynamical diffraction & non-linear imaging often “don’t matter”
- Why?
 - Kinematical is violated statistically, not systematically in most cases
 - Hence it is a valid first approximation

When does it work?

- 
- Exact problems (Diffractive Imaging)
 - Kinematical Diffraction (surfaces)
 - 1s-Channelling (HREM+HAADF)
 - Intensity ordering (PED)

L. D. Marks, W. Sinkler, Sufficient conditions for direct methods with swift electrons. *Microsc. Microanal.* **9**, 399 (2003).

Exact Cases

- 
- Suppose we have N pixels, and $N/2$ are known to be zero (compact support)
 - Wave is described by $N/2$ moduli, $N/2$ phases (for a real wave) in reciprocal space
 - Unknowns – N ; measurements $N/2$; constraints $N/2$
 - Problem is in principle fully solveable
(It can be shown to be unique in 2 or more dimensions, based upon the fundamental theorem of algebra)

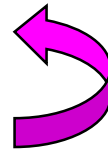
Example: Diffractive imaging

- Constraint: part of real-space x is zero
(Convex constraint)

- Iteration

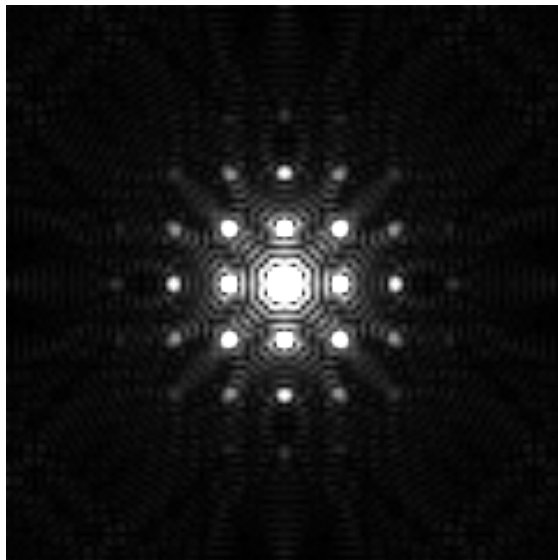
- $x = 0$, part of map

- $|X| = |X_{\text{observed}}|$



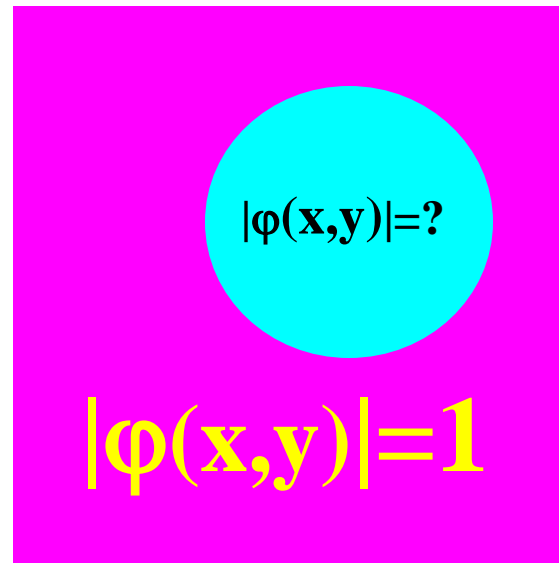
Iterate

Phase Recovery for a Small Particle



True diffraction pattern
for small particle model
(Non-Convex Constraint)

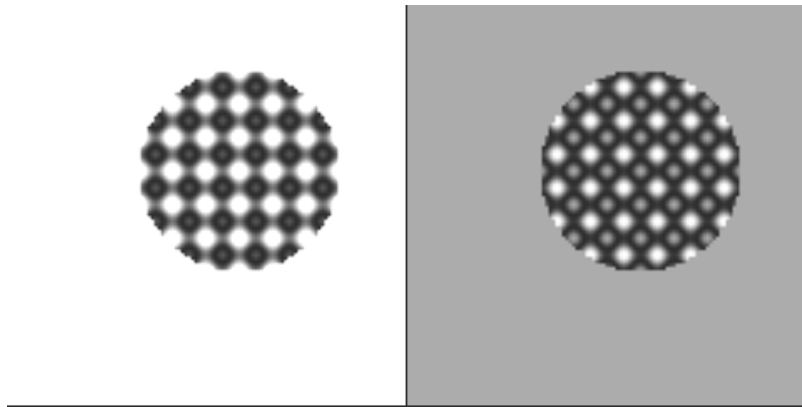
+



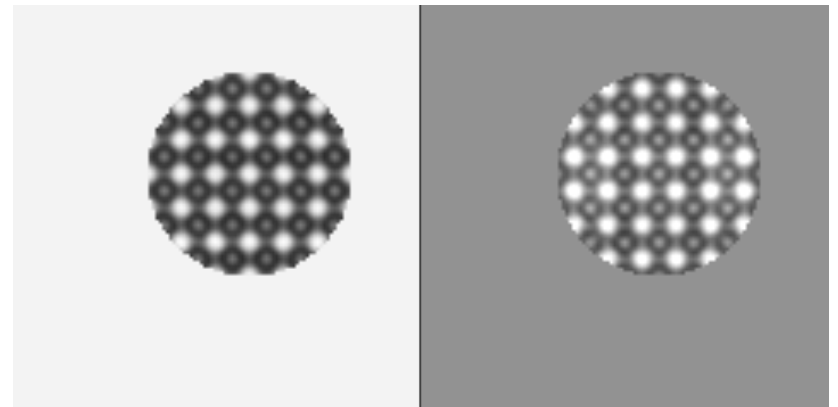
Convex Support
Constraint

= ?

Phase Recovery for a Small Particle

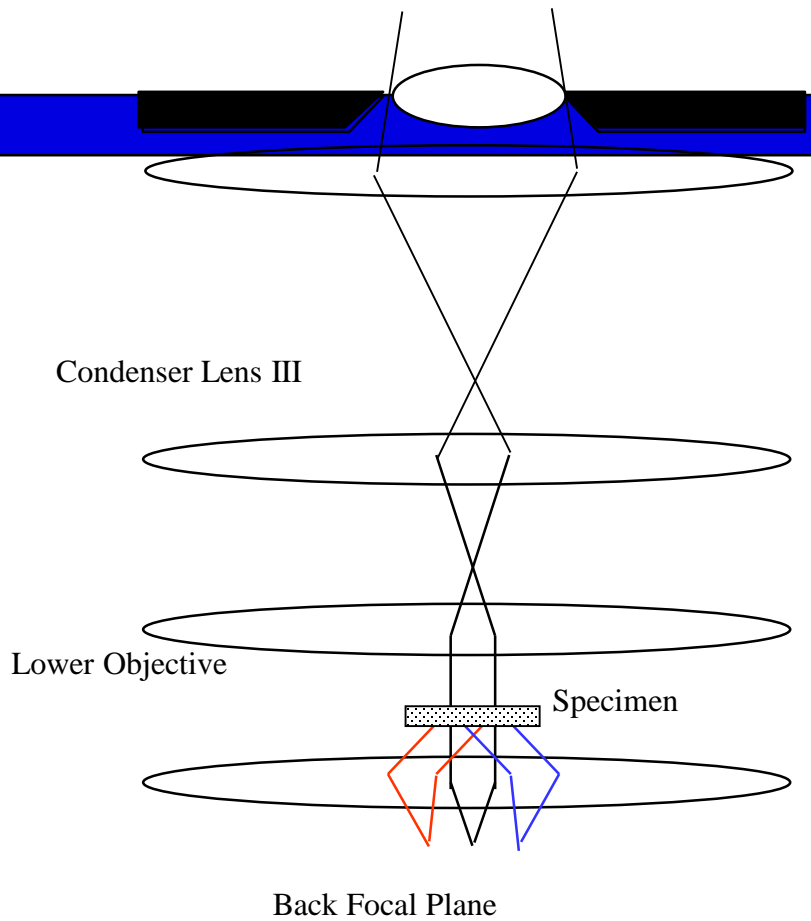


True real space exit wave for small particle model

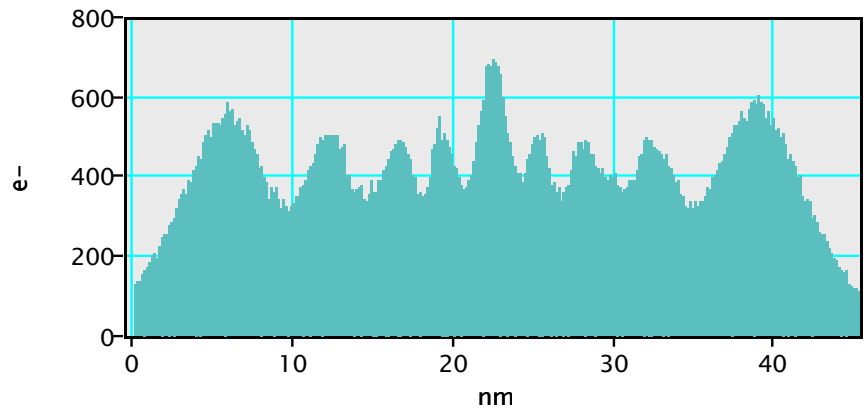
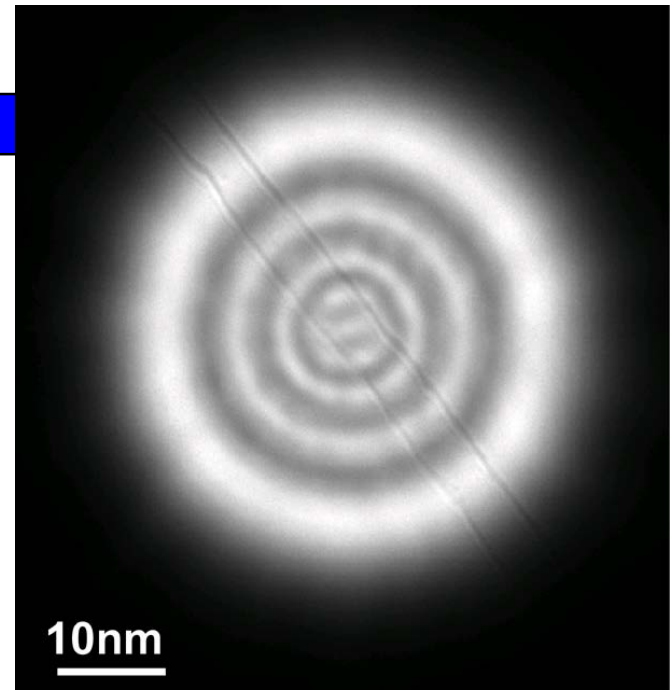


Reconstructed exit wave after 3000 iterations

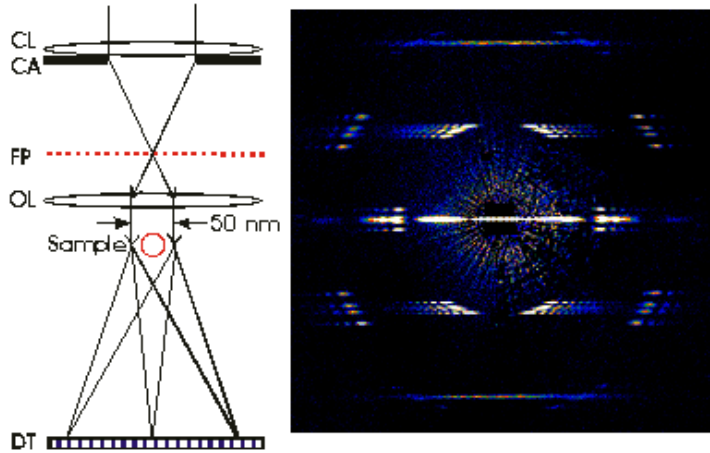
Electron Nanoprobe formation



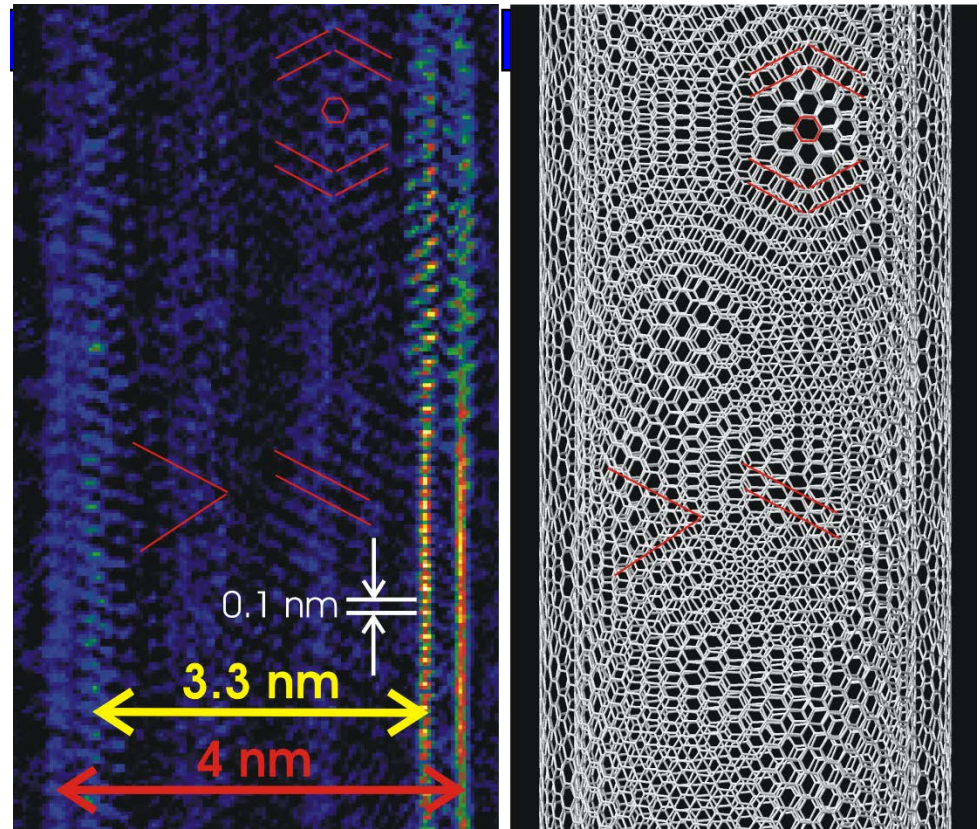
10 mm aperture \rightarrow 50 nm beam
 $M = 1/200$



Diffractive Imaging and Phase Retrieval

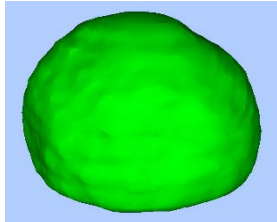
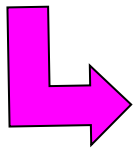
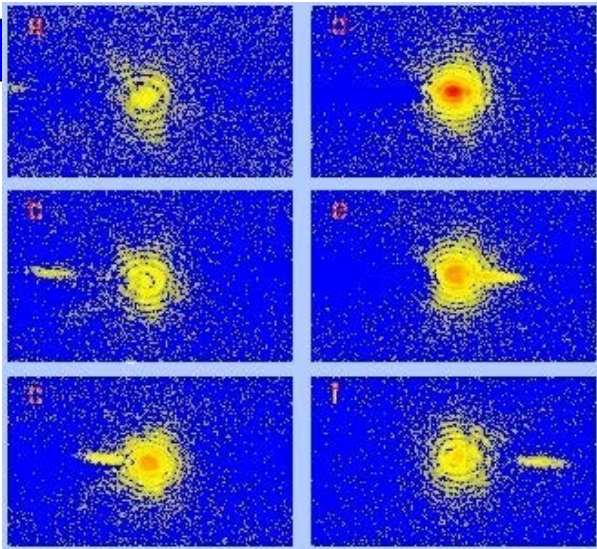


(left) A single double wall nanotube is illuminated with a narrow beam of electrons. (right) The diffraction pattern of the tube

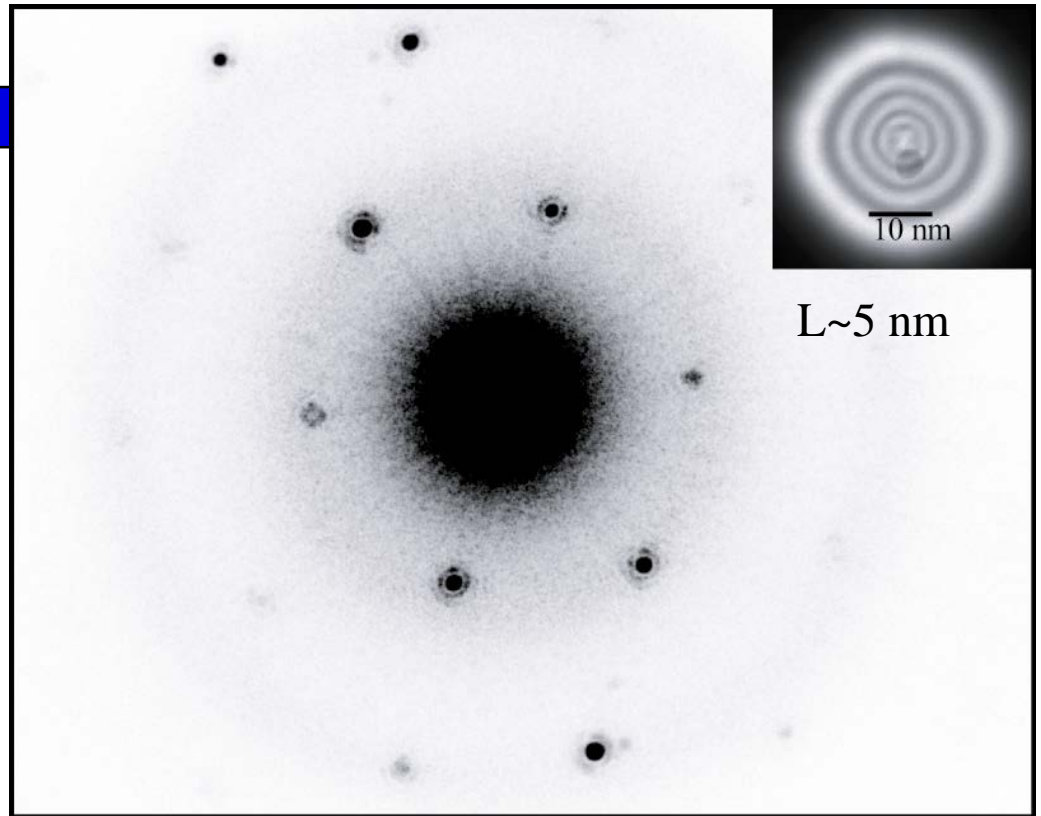


J.M. Zuo, I. Vartanyants, M. Gao, R. Zhang and L.A. Nagahara, *Science*, 300, 1419 (2003)

Single Particle Diffraction




I. Robinson
Synchrotron



- Atomic resolution
- Strong interaction of electrons

J. Tao, See Zuo et al, Microscopy Research Techniques, 2004

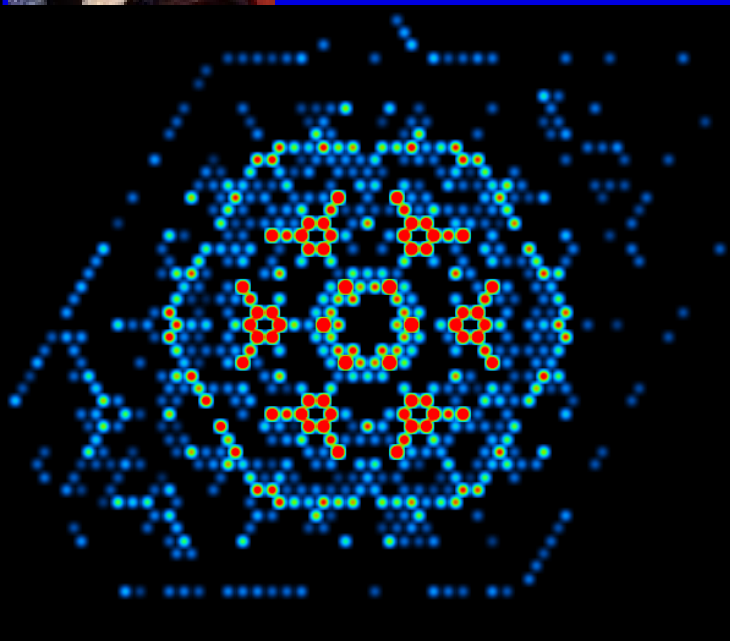
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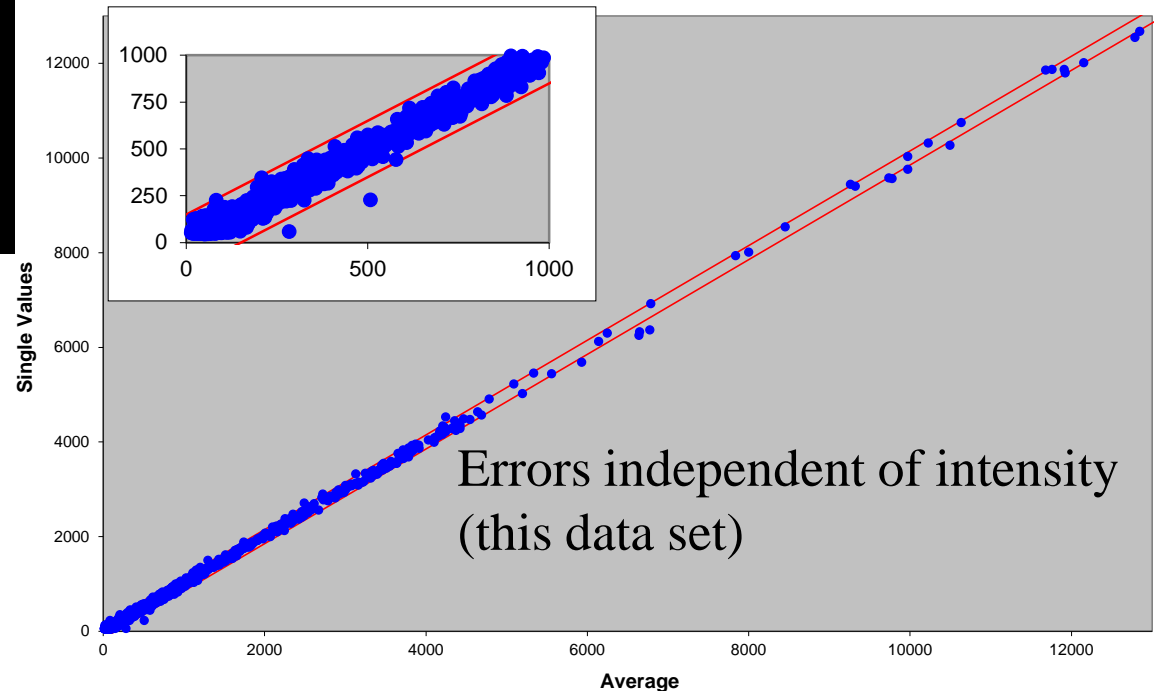
L. D. Marks, W. Sinkler, Sufficient conditions for direct methods with swift electrons. *Microsc. Microanal.* **9**, 399 (2003).



TED: Si (111) 7x7



Method: Merge data for 6-20 different exposures to obtain accuracies of ~1% with statistical significance

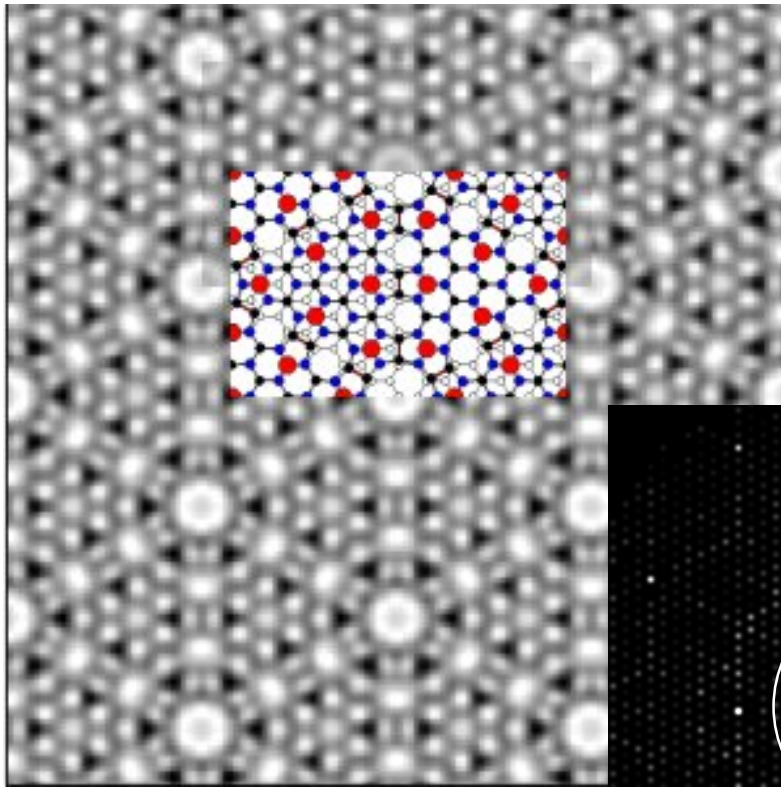


Cross-Correlation
Method

P. Xu, et al.

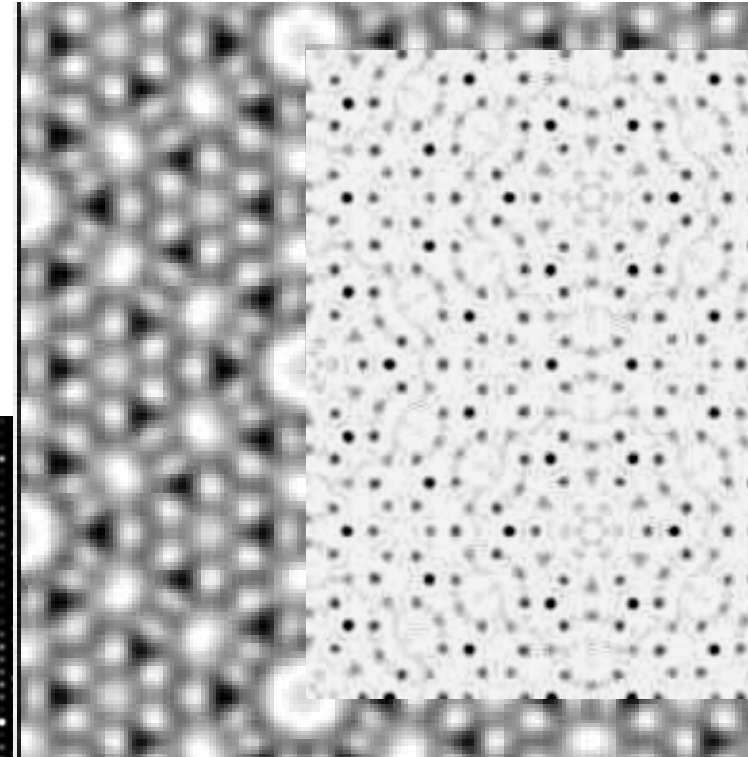
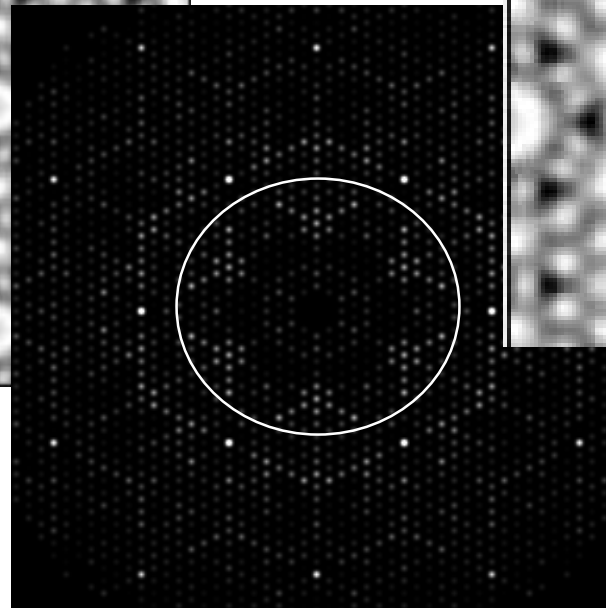
Ultramicroscopy **53**, 15
(1994).

Restoration and Extension



0.3nm Image

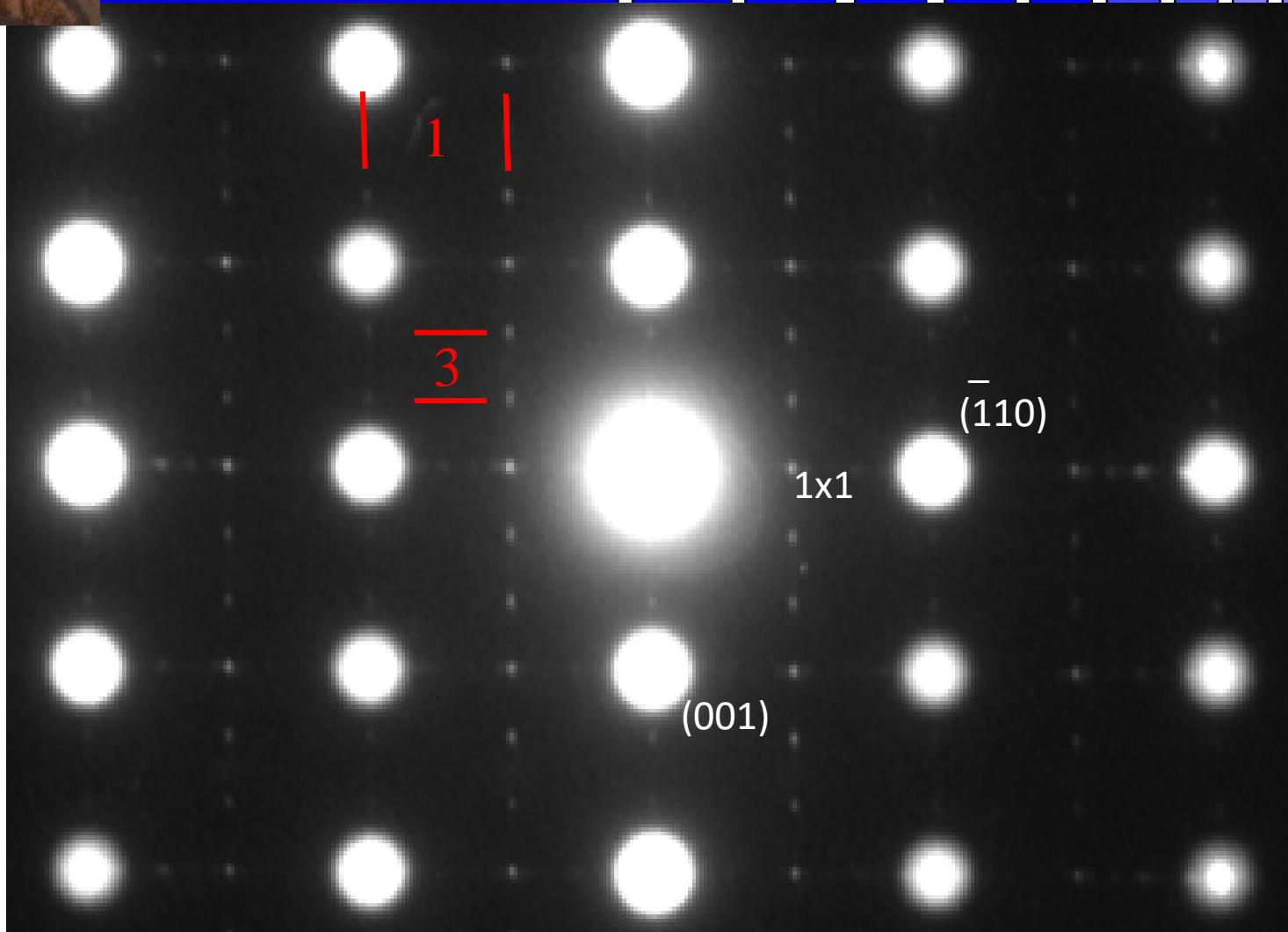
+DP



0.05nm Image

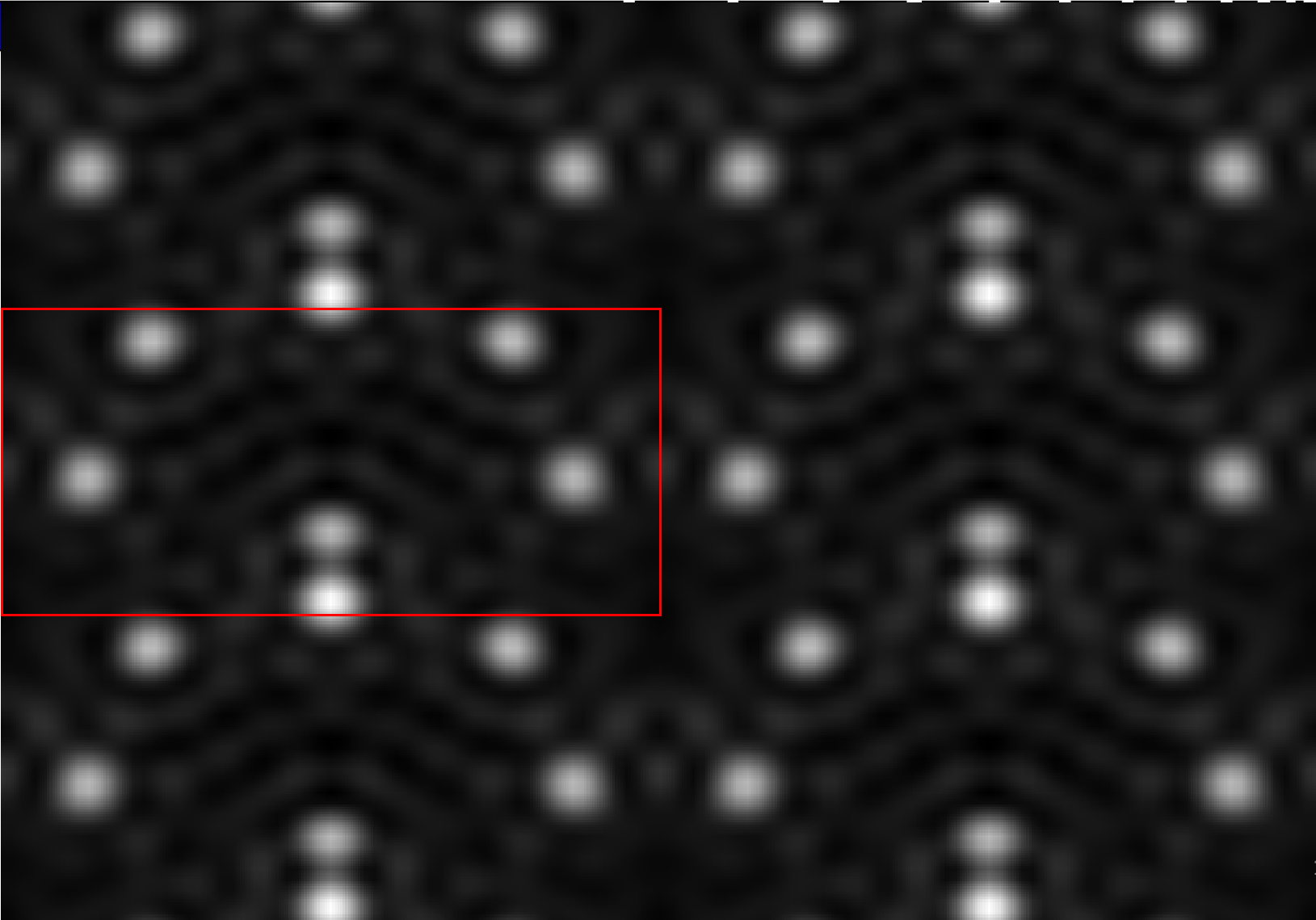


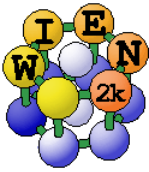
1000 °C in flowing O₂



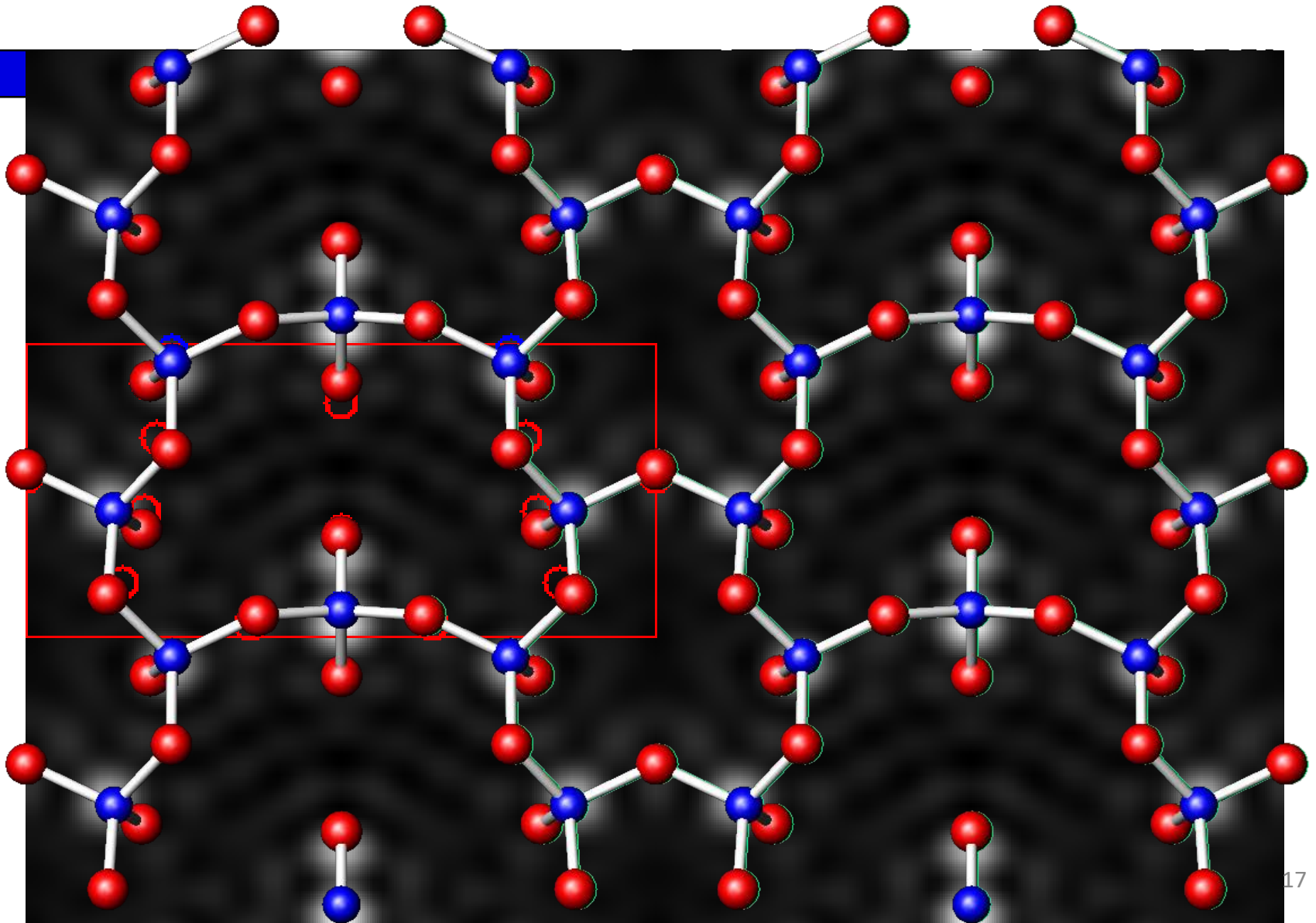


Direct Methods Solution



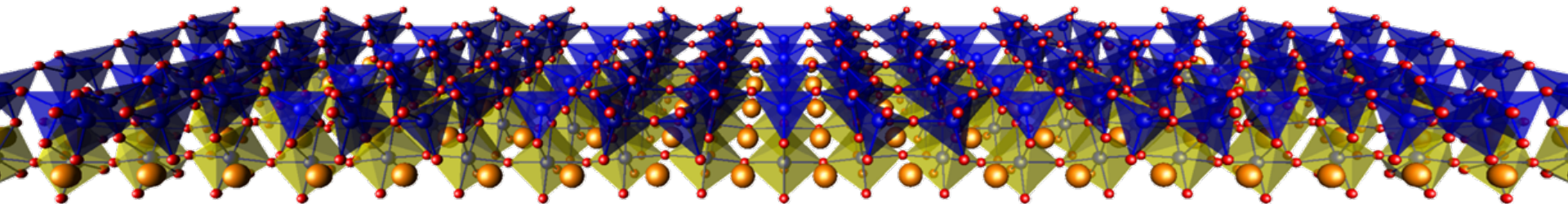
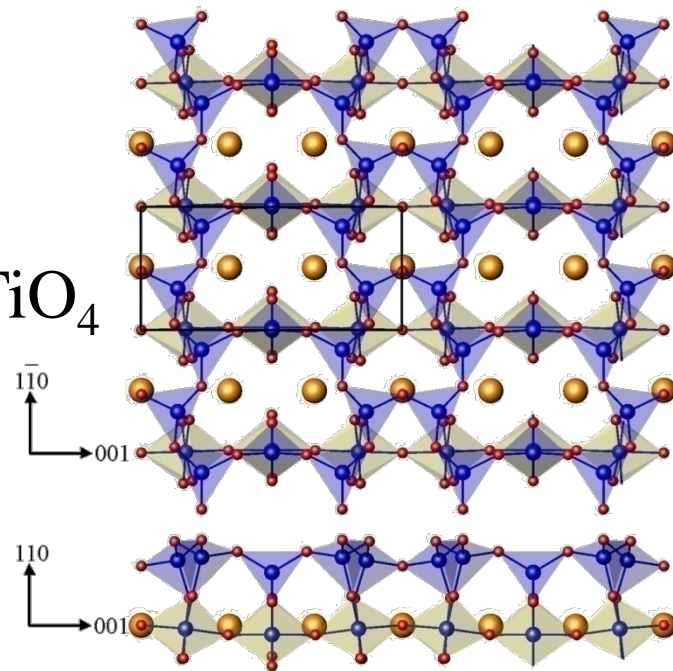


Atomic Positions Refined



SrTiO₃ (110) 3x1

- TiO₂ overall surface stoichiometry
 - Ti₅O₇ atop O₂ termination
 - Ti₅O₁₃ atop SrTiO termination
- Surface composed of corner sharing TiO₄ tetrahedra
 - Arranged in rings of 6 or 8 tetrahedra
 - 4 corner share with bulk octahedra
 - 1 edge shares with bulk octahedron



Blue polyhedra are surface polyhedra, gold are bulk octahedra, orange spheres Sr, blue spheres Ti, red spheres O

Inversion

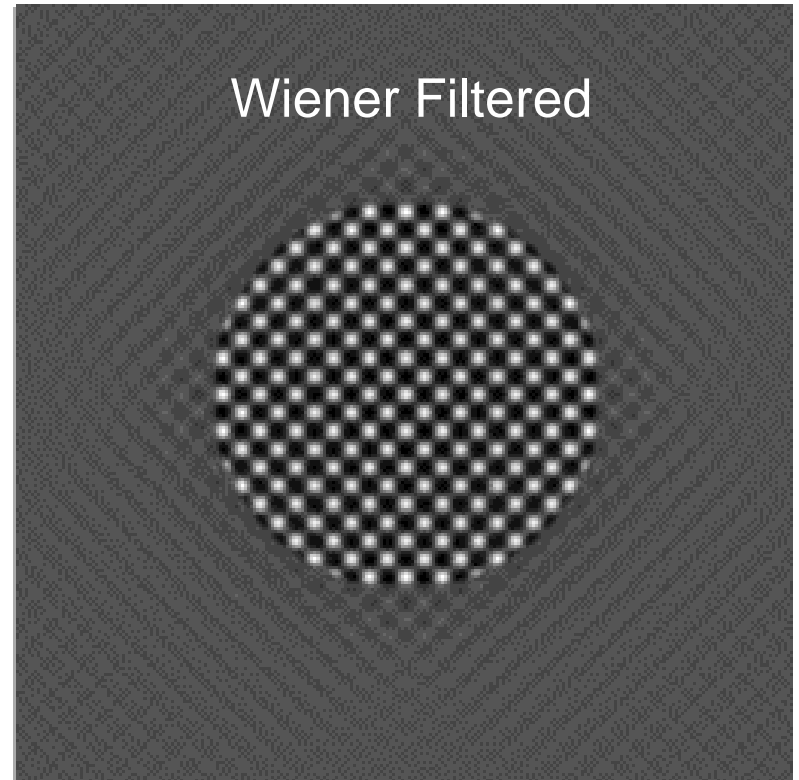
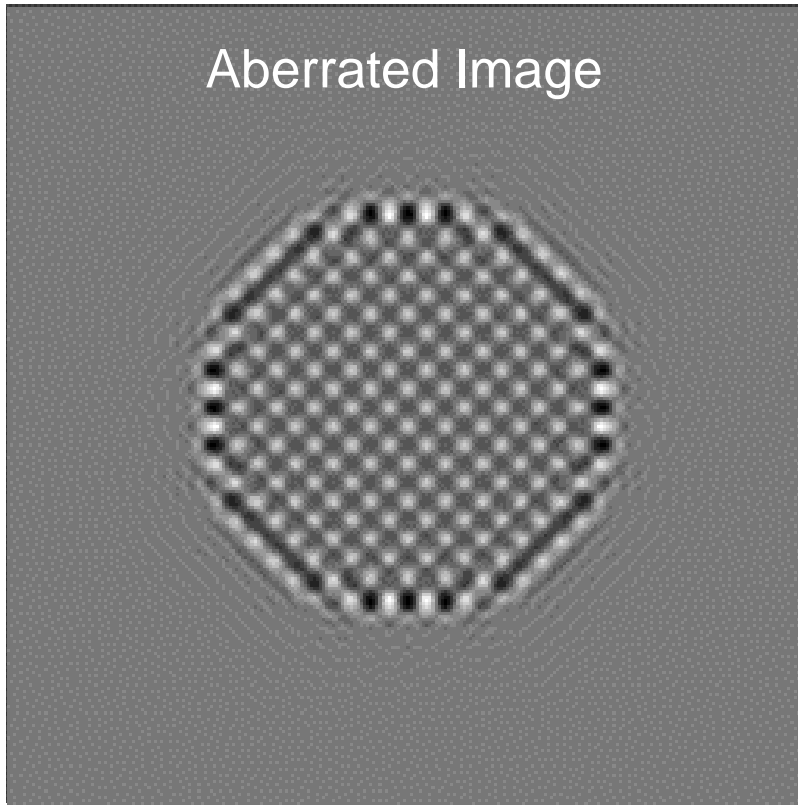
- $I(r) \sim \int \Psi(u)T(u)\exp(2\pi iu.r)du + \text{noise}$
write $A(u)=\Psi(u)T(u)$
- The optimal filter (L2) $F(u)$ to apply is given by (Wiener, 1940)

$$F(u) = T^*(u)/\{|T(u)|^2+n(u)^2/S(u)^2\}$$

$n(u)$ = spectral distribution of noise

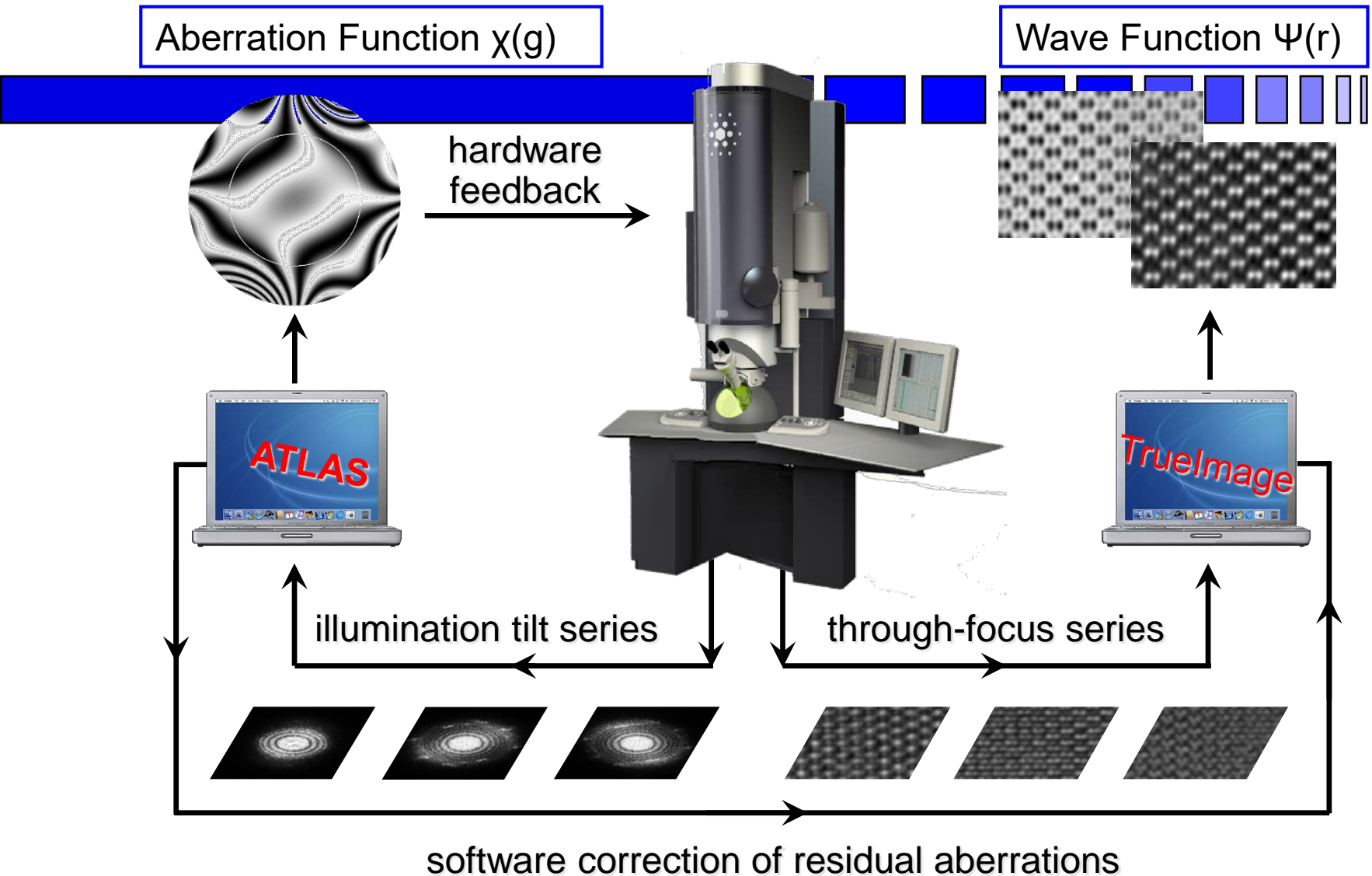
$S(u)$ = estimate of signal

Wiener Filtering

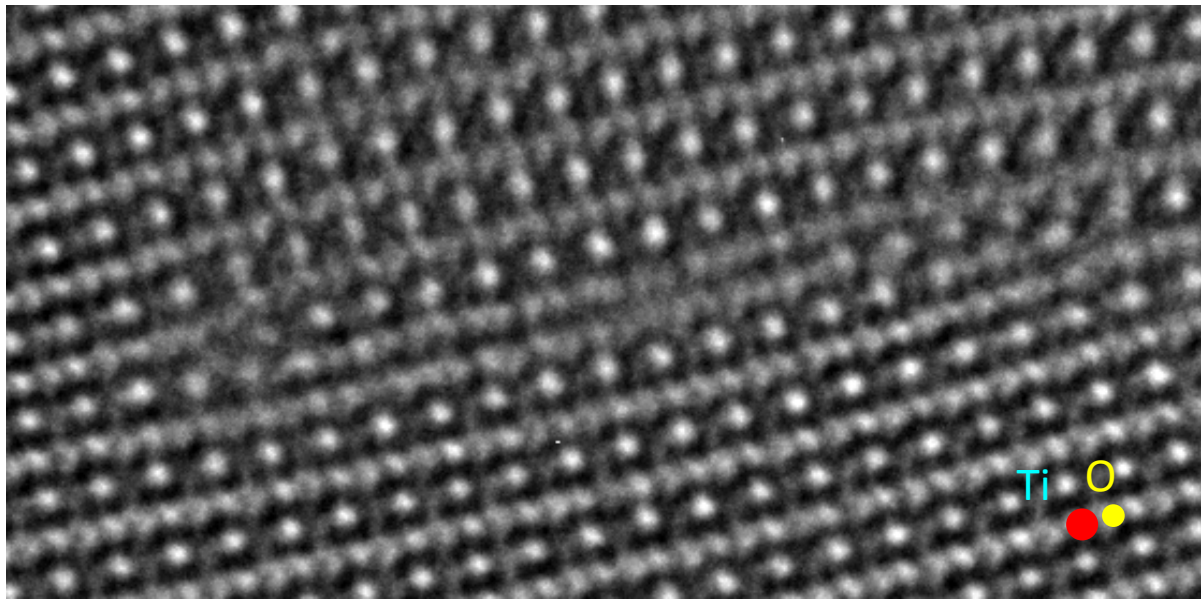


Simple Example

Aberration control & reconstruction of electron wave function



ATLAS & TruelImage:: Stacking Faults in SrTiO₃ (110)



Z_{opt} micrograph

Titan 80-300

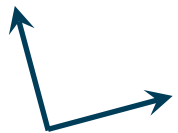
deficiencies:

shaded columns

inf. signal-to-noise ratio

spurious contrast peaks

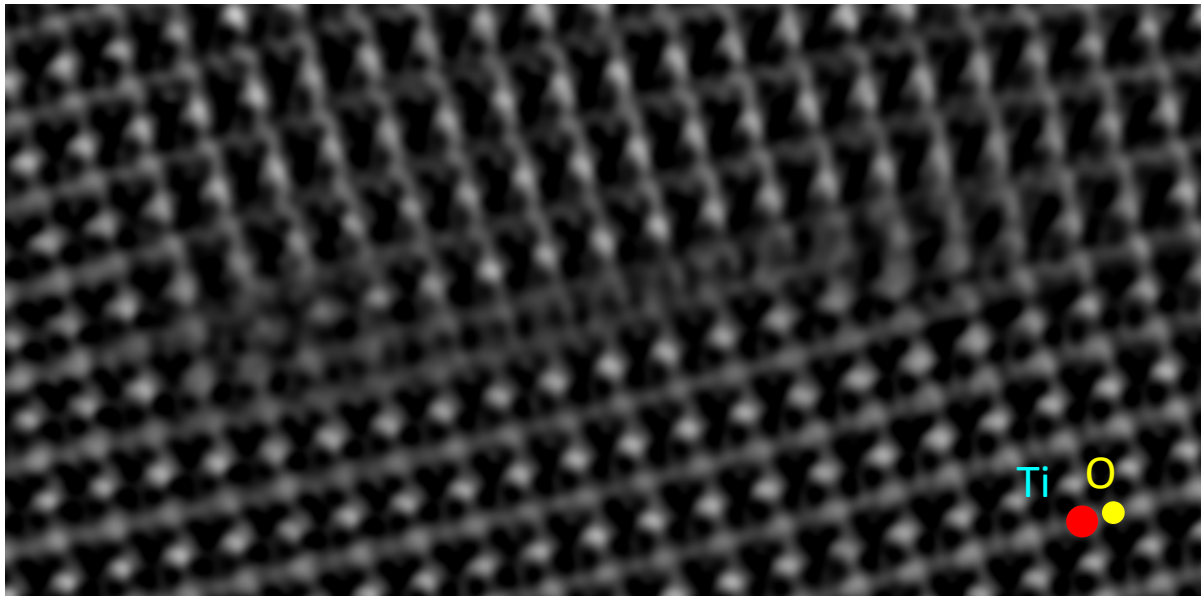
[001]



[110]

1.38 Å

ATLAS & TruelImage:: Stacking Faults in SrTiO₃ (110)



uncorrected phase image

Titan 80-300

deficiencies:

shaded columns

[001]

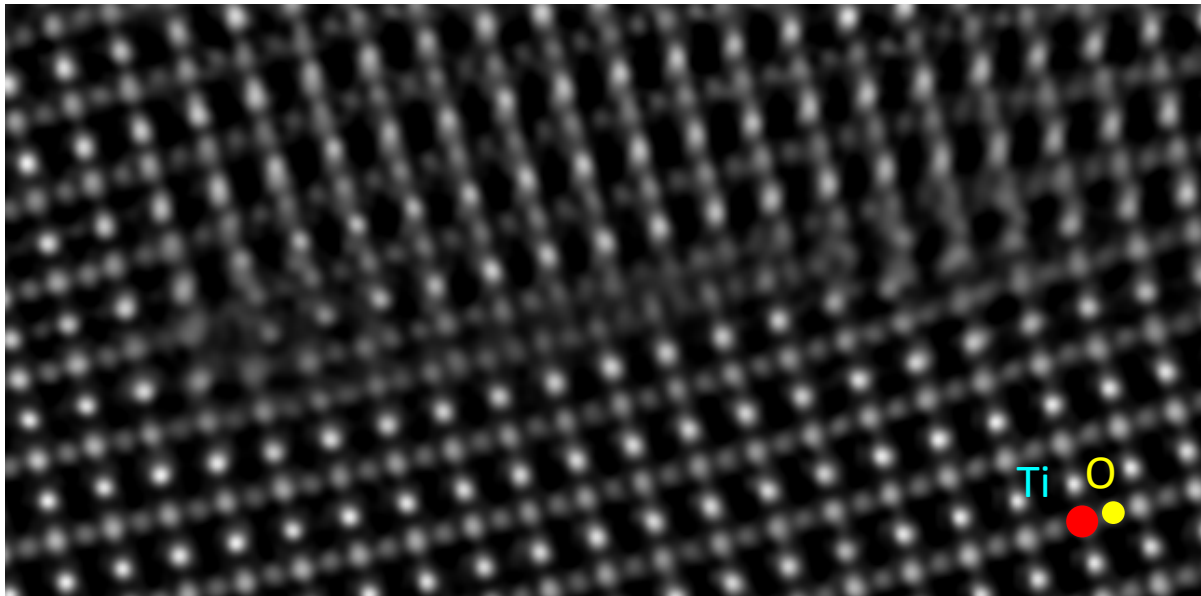


[$\bar{1}10$]



1.38 Å

ATLAS & TruelImage:: Stacking Faults in SrTiO₃ (110)



corrected phase image

Titan 80-300

deficiencies:

none

[001]




[$\bar{1}10$]

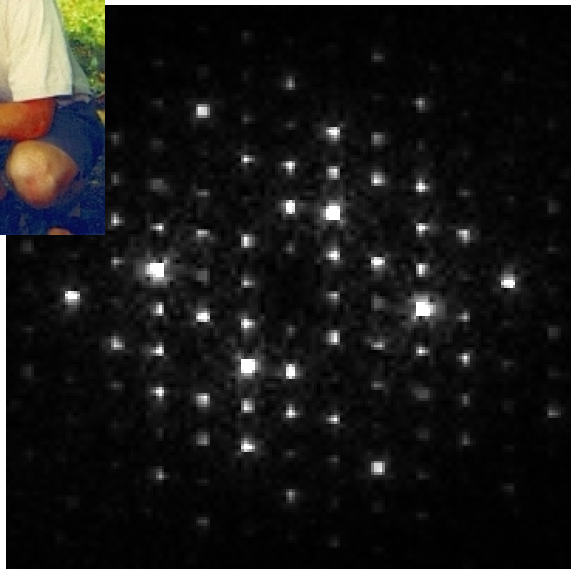


1.38 Å

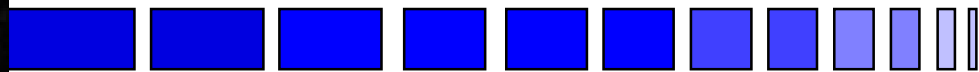
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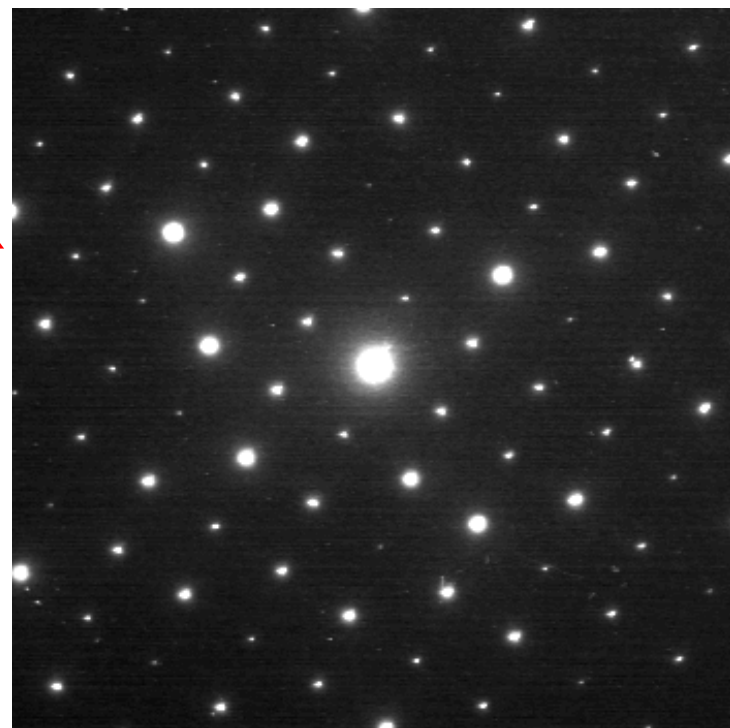


FFT

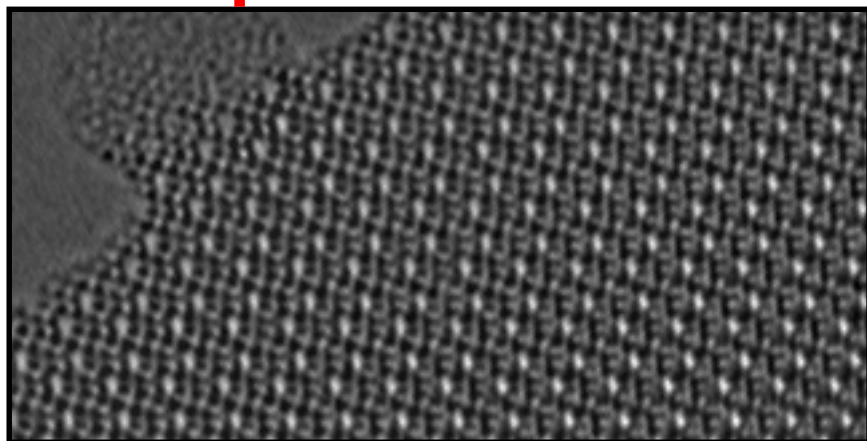


Method

Initial
Phases

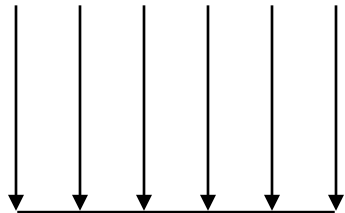


Nanoprobe
Diffraction Data

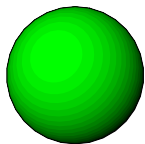
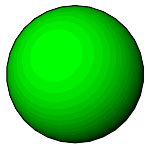
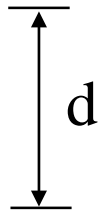
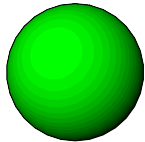
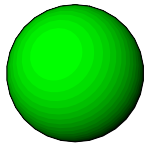


Conventional HREM Image

Channeling Approximation



$$e^- \quad V(\mathbf{x}, \mathbf{y}) = \frac{1}{d} \int_{-\infty}^{\infty} V_0(\mathbf{r}) dz$$

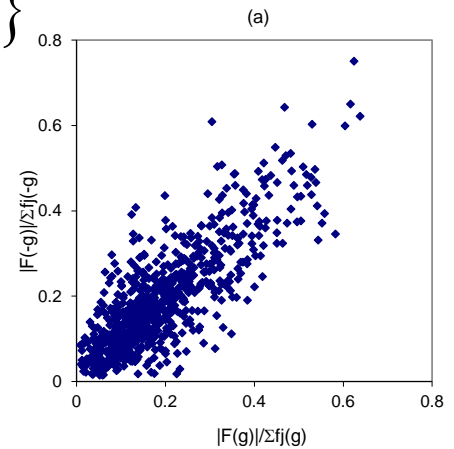


$$\psi(\mathbf{r}, \mathbf{z}) = \sum_n \mathbf{C}_n \Phi_n(\mathbf{r}) \exp\{-i\lambda_n \mathbf{z}\}$$

λ_n 2-D Eigenvalue

Talks later by Van Dyck and Chukhovskii will explain more details

Σ_0 distribution is statistically kinematical



F. N. Chukhovskii, et al *Acta Cryst A* **57**, 231 (2001)

Statistics in a 1s model



Kinematical

Σ_0 :

$$\phi(\mathbf{g}) + \phi(-\mathbf{g}) = 0$$

$$|\mathbf{F}(\mathbf{g})| = |\mathbf{F}(-\mathbf{g})|$$

Σ_2 :

$$\phi(\mathbf{g}) + \phi(\mathbf{h}-\mathbf{g}) + \phi(-\mathbf{h}) \sim 2n\pi$$

Dynamical

Σ_0 :

$$\phi(\mathbf{g}) + \phi(-\mathbf{g}) \sim \omega$$

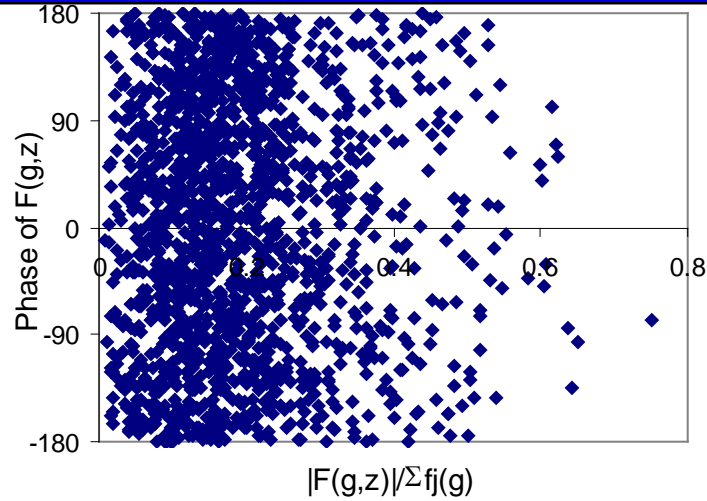
$$|\mathbf{F}(\mathbf{g})| \sim |\mathbf{F}(-\mathbf{g})|$$

Σ_2 :

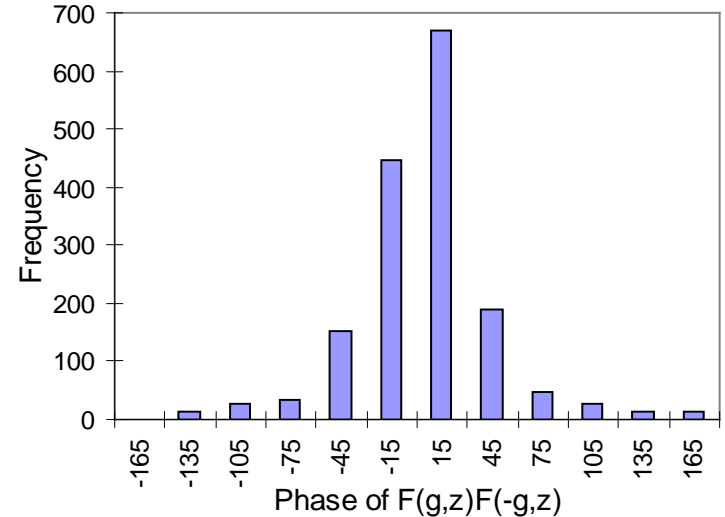
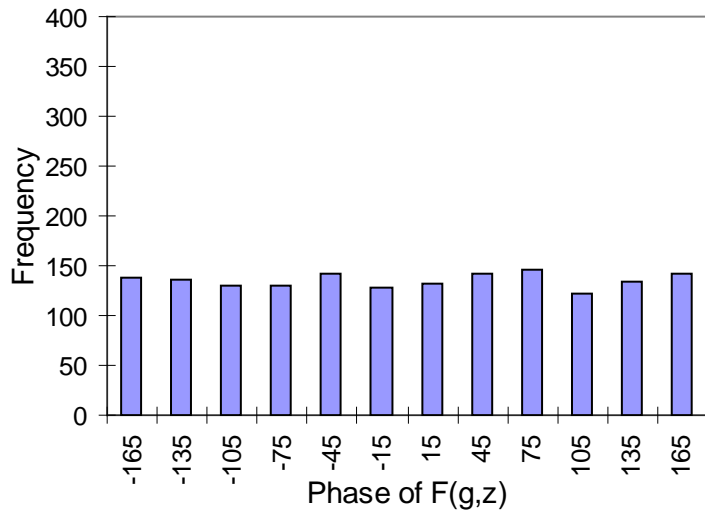
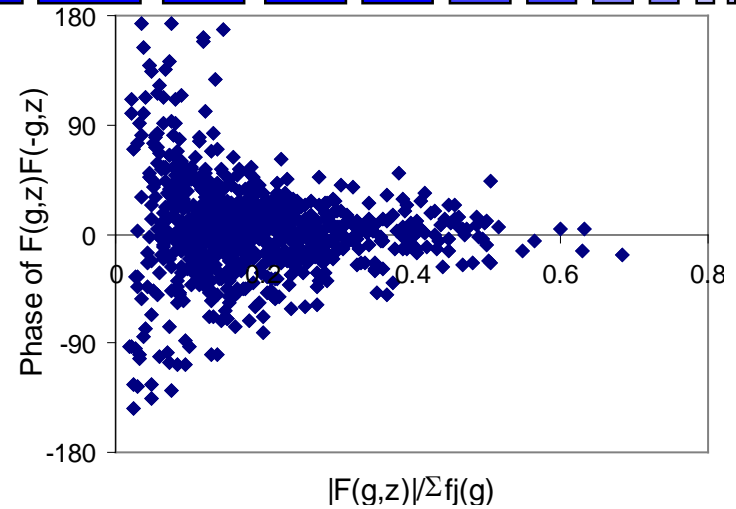
$$\phi(\mathbf{g}) + \phi(\mathbf{h}-\mathbf{g}) + \phi(-\mathbf{h}) \sim \zeta$$

Kinematical Theory for electrons is “wrong”, but statistically “right” in many cases, **BUT NOT ALL** (see C252)

Σ_0 dynamical distribution

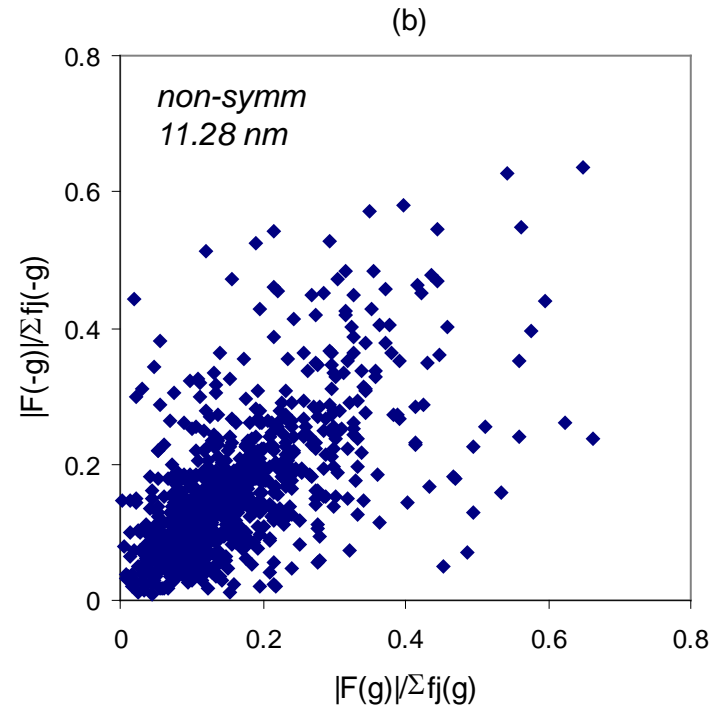
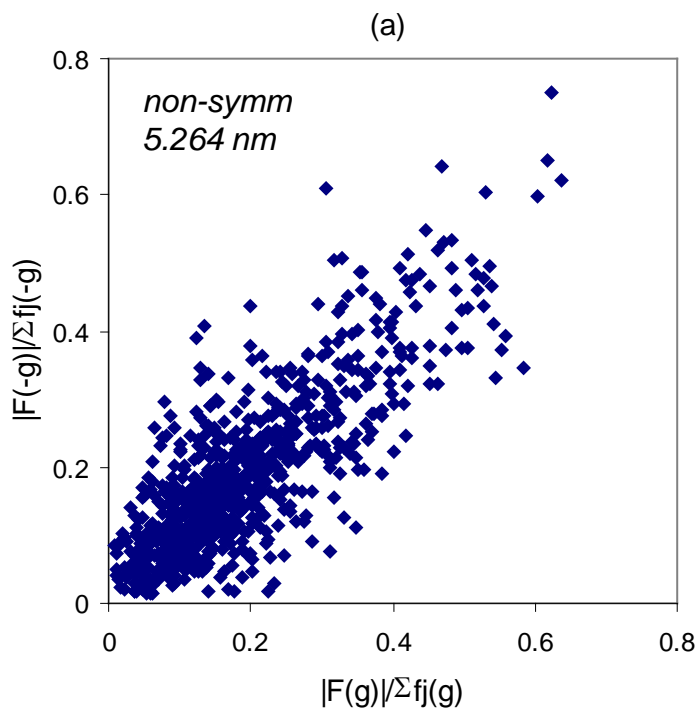


non-symm
5.264nm

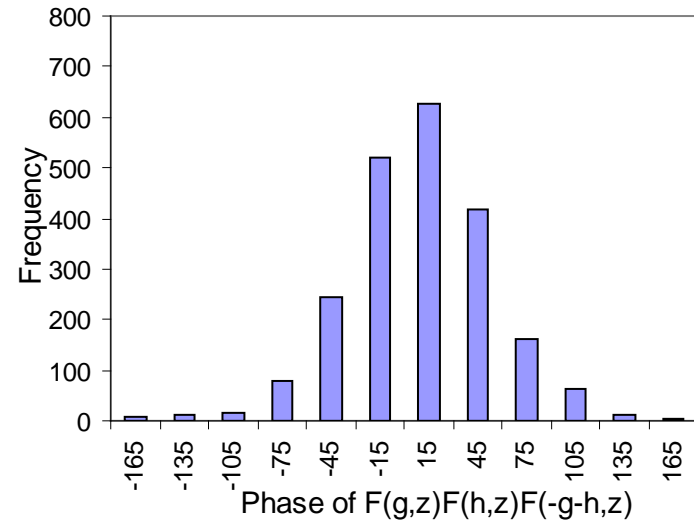
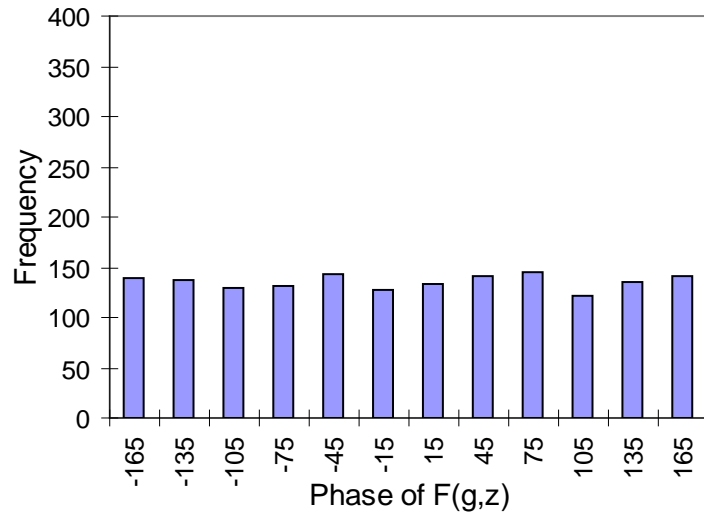
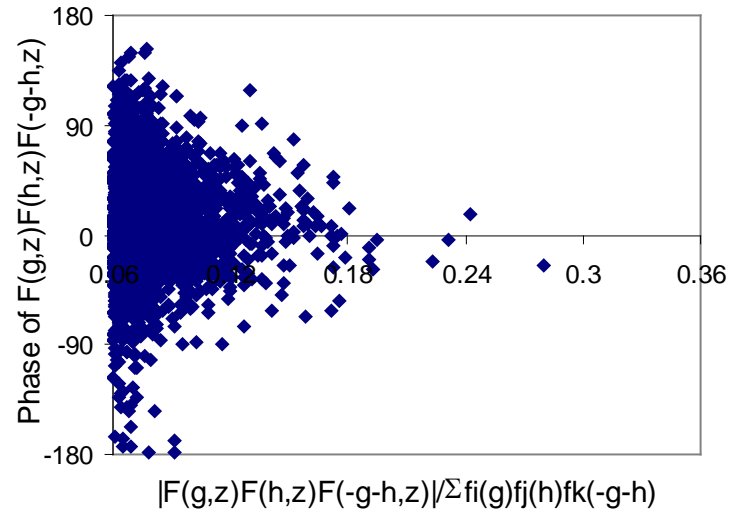
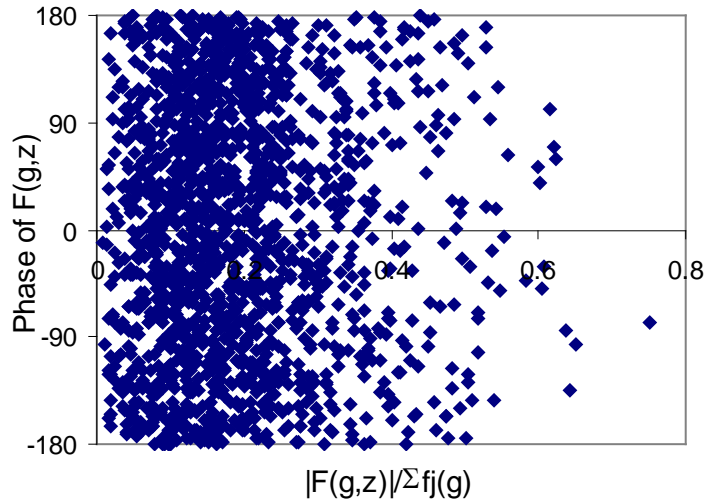


Σ_0 dynamical distribution

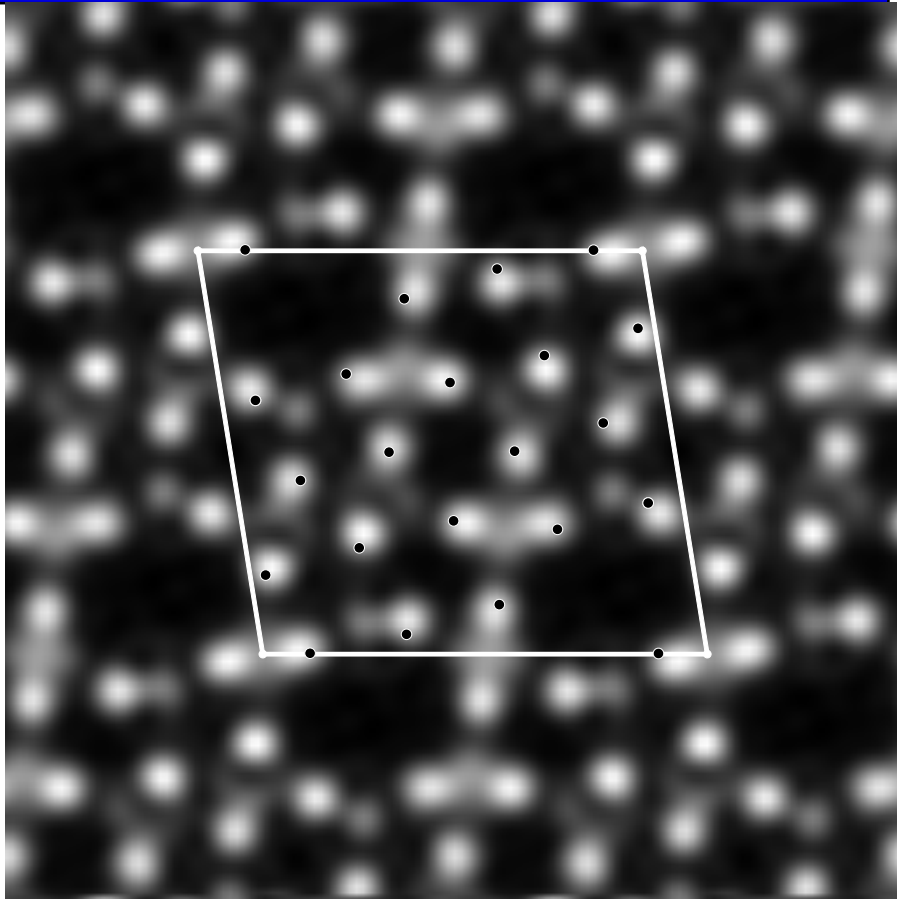
Freidel symmetry is statistical



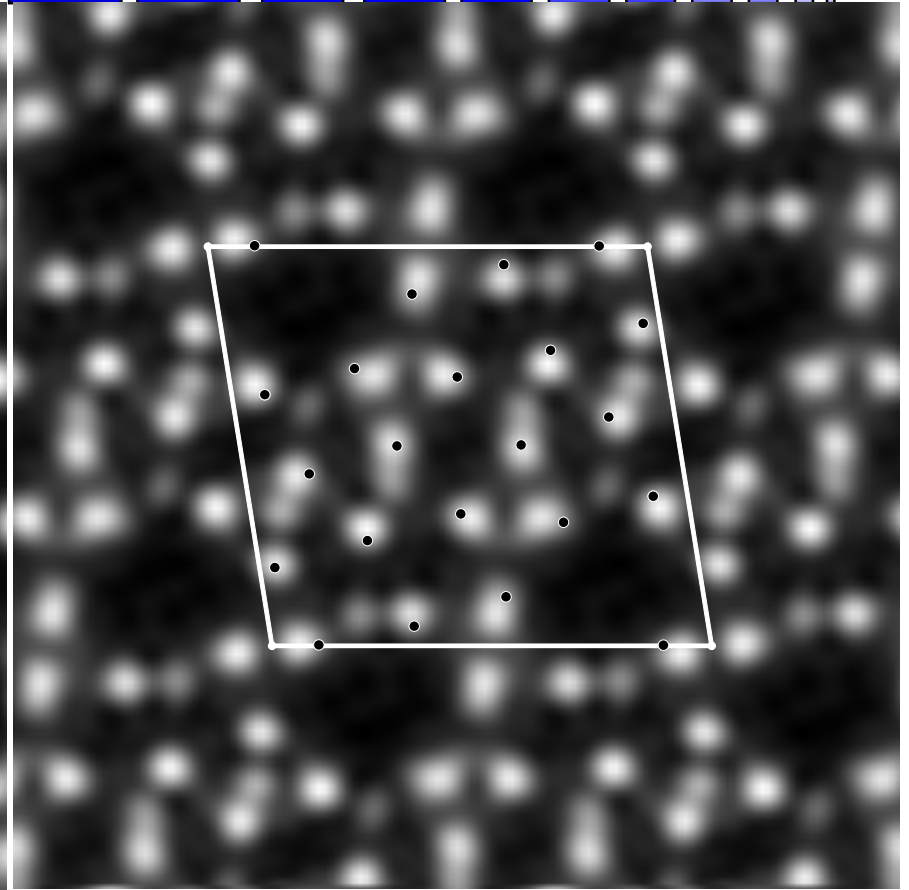
Σ_2 dynamical distribution



Calculated Wave

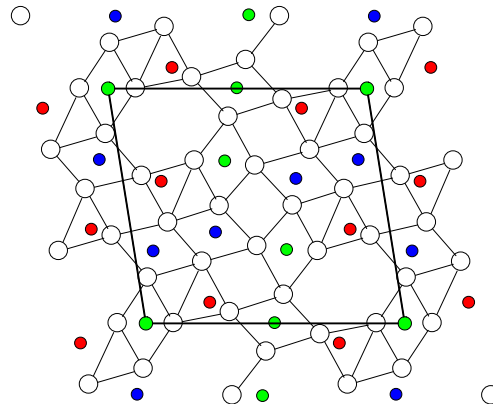
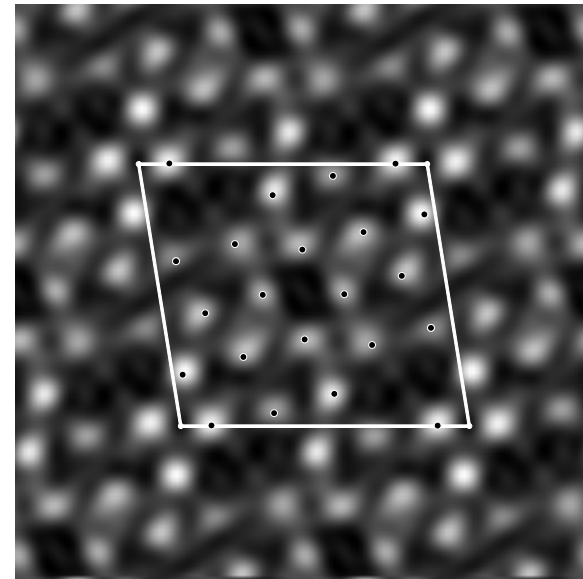
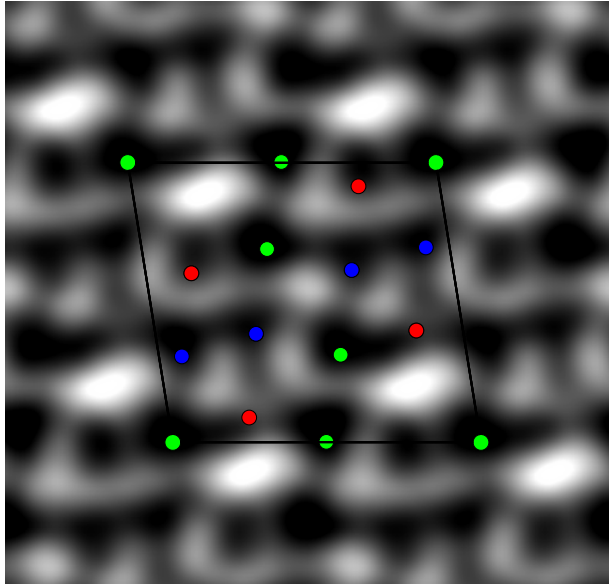


$|\psi(\mathbf{r})-1|$ at 113 Å thickness




$|\psi(\mathbf{r})-1|$ at 202 Å thickness

O sites in $(\text{Ga,In})_2\text{SnO}_5$ determined using direct phasing of TED data.




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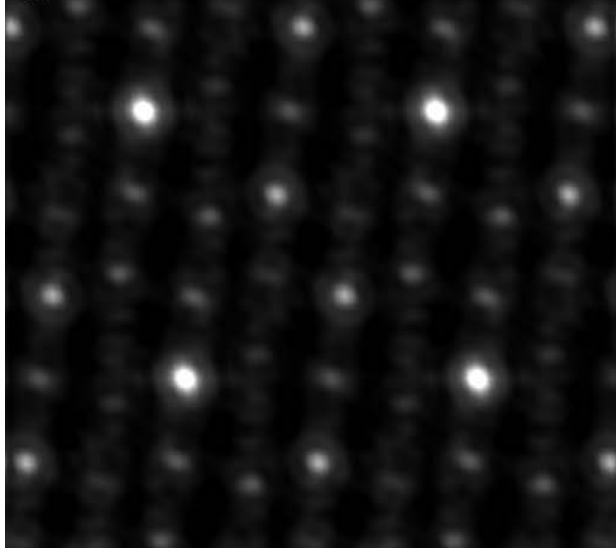
Precession Electron Diffraction

- 
- Quasi-Kinematical Data
 - Averaging over angle/phase (and thickness) damps dynamical contributions
 - Intensities are close to monatomic with structure factors (statistically)



(Ga,In)₂SnO₄ precession data: Direct methods solution

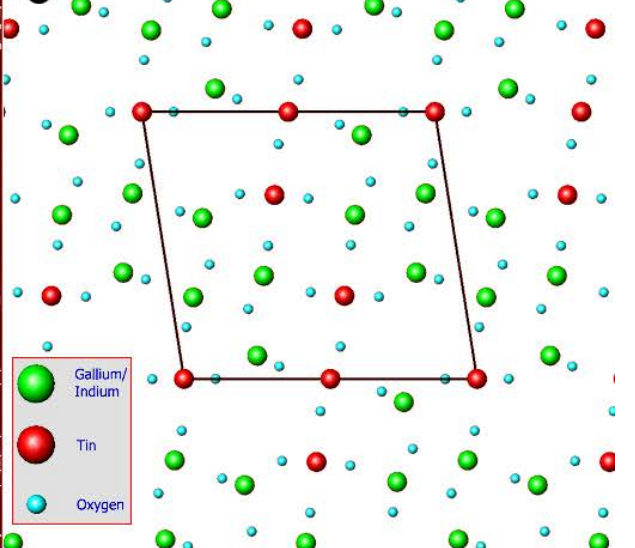
a (Real Space)



b



c



	ΔR (Å)
Sn1	0.00E+00
Sn2	0.00E+00
Sn3	6.55E-03
In/Ga1	5.17E-02
In/Ga2	2.37E-03
Ga1	6.85E-02
Ga2	1.22E-01


Displacement ($R_{\text{neutron}} - R_{\text{precession}}$):

$$\Delta R_{\text{mean}} < 4 \cdot 10^{-2} \text{ \AA}$$

(Sinkler, et al. J. Solid State Chem, 1998).

(Own, Sinkler, & Marks, submitted.)

Conclusion

- 
- The “Phase Problem” with electrons is no longer really a problem....assuming ideal data of course
 - Many techniques work most of the time
 - Few techniques work all the time
 - Some unresolved issues (proper dynamical refinement)
 - Remember that we are solving an inversion problem, and these are susceptible to ill-conditioning

Four basic elements are required to solve a recovery problem



1. A data formation model

Imaging/Diffraction/Measurement

2. A priori information

The presence of atoms or similar

3. A recovery criterion:

A numerical test of Goodness-of-Fit

4. A solution method.

Mathematical details